

Light Stop Decays

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in collaboration with Ramona Gröber, Margarete Mühlleitner and Eva Popena, arXiv:1408.4662

INSTITUTE FOR THEORETICAL PHYSICS



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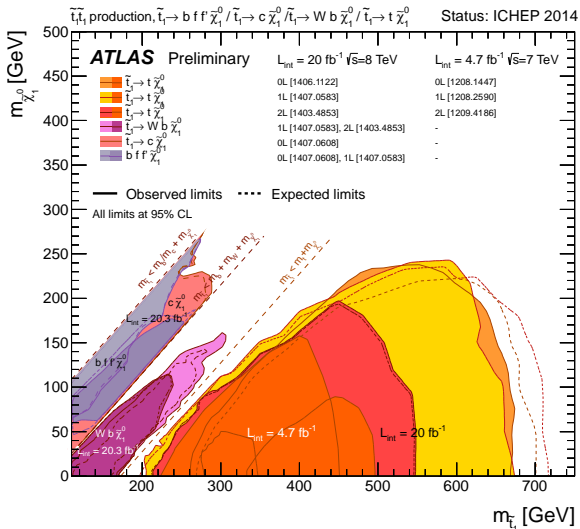
Supersymmetry (SUSY)

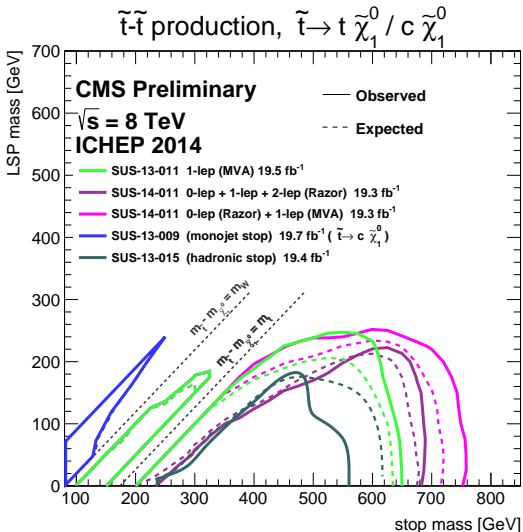
- extension of the Standard Model (SM)
→ Dark Matter, gauge coupling unification, stabilization of Higgs mass
- SM gauge group
- enhanced space-time symmetry
- enriched particle spectrum → superpartners (s..., ...ino, $\tilde{}$), extended Higgs sector
- conserved SUSY → masses of superpartners = masses of SM particles
- soft SUSY breaking terms: general parametrization
- R-parity → lightest SUSY particle (LSP) stable, production of SUSY particles in pairs, decays of SUSY particles have odd numbers of SUSY particles in the final state

Light Stops and Their Decays

- large mass splitting possible for stops
- light stops compatible with experimental results
- stops accessible at LHC energies
- $\Delta m = m_{\tilde{t}_1} - m_{\tilde{\chi}_1^0}$; $\Delta m < m_W$
- existing work:
 - $\tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ electroweak one-loop process
[Hikasa, Kobayashi '87][Jahn '98][Mühlleitner, Popenda '11]
 - four-body decay without FV
[Boehm, Djouadi, Mambrini '99]
- here:
 - FCNC decay $\tilde{u}_1 \rightarrow (c, u)\tilde{\chi}_1^0$ at NLO SUSY-QCD
 - four-body decay $\tilde{u}_1 \rightarrow \tilde{\chi}_1^0 d_j \bar{f} \bar{f}'$ with final state mass effects
- general flavor structure
- MSSM; LSP: $\tilde{\chi}_1^0$; NLSP: \tilde{u}_1

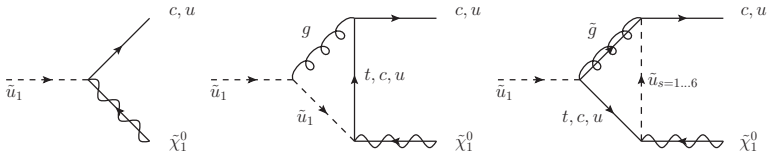
Introduction: Experimental Status



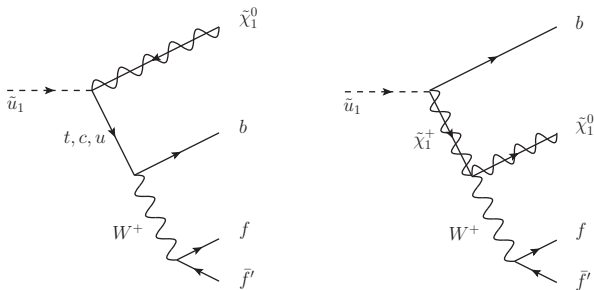


Decays - Examples of Feynman Diagrams

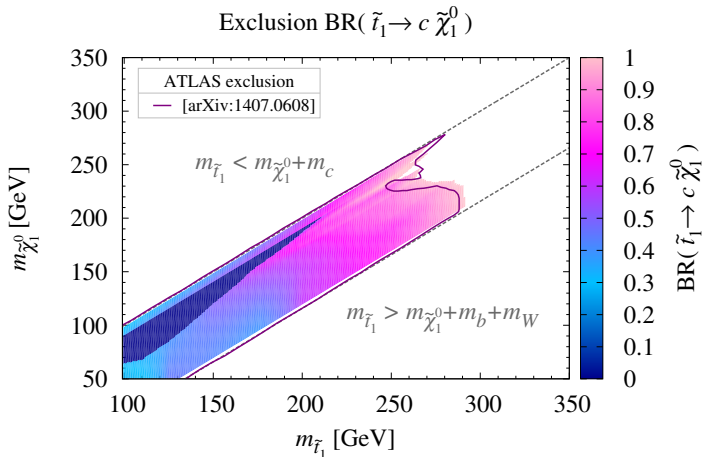
Examples for the Two-Body Decay

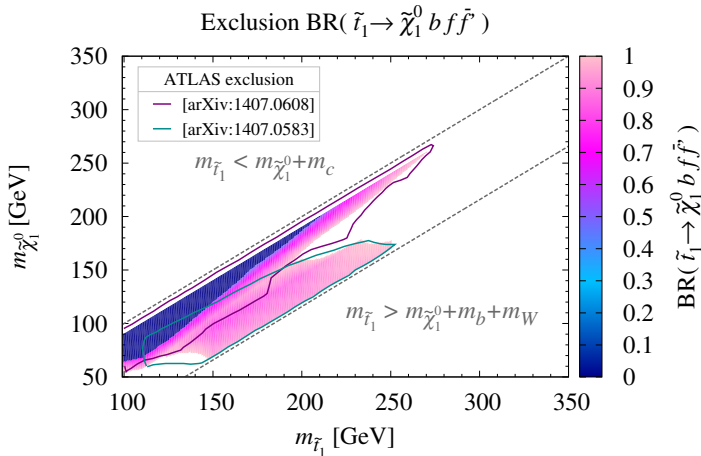


Examples for the Four-Body Decay



- implementation of the decays: `SUSYHIT` [Djouadi, Mühlleitner, Spira '12]
[Gröber Mühlleitner, Popenda, AW '14]
- spectrum generator: `SPheno` [Porod '12]
- Higgs decays, BRs, effective couplings: `HDECAY` [Djouadi, Kalinowski, Mühlleitner, Spira '10]
- Higgs observation and exclusion bounds: `HiggsBounds` and `HiggsSignals`
[Bechtle, Brein, Heinemeyer, Stål, Stefaniak, Weiglein, Williams '13]
- Relic Density, $\Omega_c h^2 < 0.12$: `SuperIso Relic` [Arbey, Mahmoudi '11]
- B-physics observables: `SuperIso` [Mahmoudi '09]
 $BR(b \rightarrow X_s \gamma)$, $BR(B \rightarrow \tau \nu)$, $BR(B_{(s)}^0 \rightarrow \mu^+ \mu^-)$
- constraints on relevant SUSY masses and on the lightest CP-even Higgs boson mass
 $m_{\tilde{g}} > 1.45 \text{ TeV}$ [CMS-SUS-13-007, ATL-PHYS-PROC-2013-179]
 $m_{H^0} = (125.5 \pm 3.0) \text{ GeV}$ [CMS-HIG-12-028, CERN-PH-EP-2012-218]
 $m_{\tilde{u}_1}$, $m_{\tilde{\chi}_1^0}$, \rightarrow exclusion limits

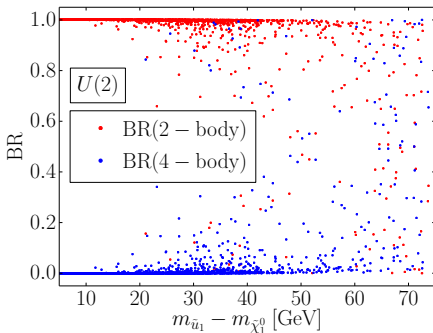
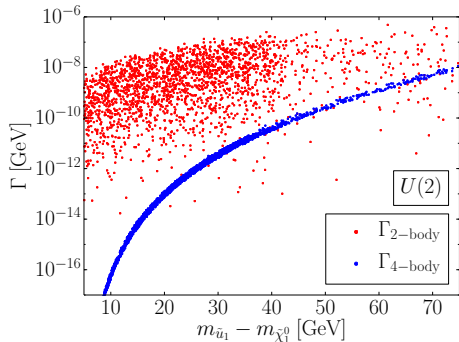




Results: Random Scan over

$A_t, M_1, m_{\tilde{U}_3}, m_{\tilde{Q}_{L3}}, \tan \beta, m_A$

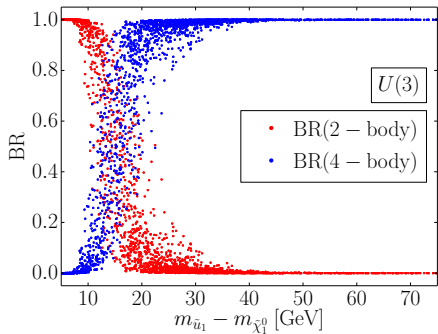
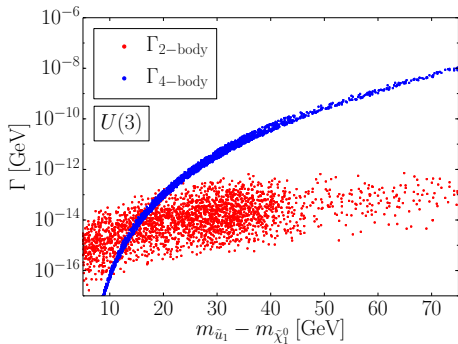
$$m_{\tilde{Q}_{L1}} = m_{\tilde{Q}_{L2}} \neq m_{\tilde{Q}_{L3}}$$



Results: Random Scan over

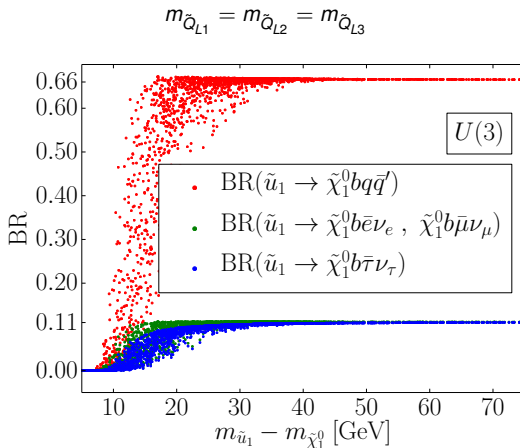
$A_t, M_1, m_{\tilde{U}_3}, m_{\tilde{Q}_L}, \tan \beta, m_A$

$$m_{\tilde{Q}_{L1}} = m_{\tilde{Q}_{L2}} = m_{\tilde{Q}_{L3}}$$

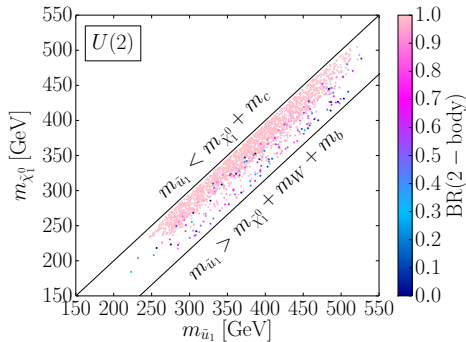


Results: Random Scan over

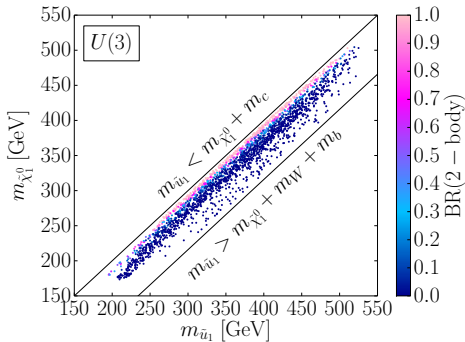
$A_t, M_1, m_{\tilde{U}_3}, m_{\tilde{Q}_L}, \tan \beta, m_A$



Results: Direct Comparison



Random scan with
 $m_{\tilde{Q}_{L1}} = m_{\tilde{Q}_{L2}} \neq m_{\tilde{Q}_{L3}}$



Random scan with
 $m_{\tilde{Q}_{L1}} = m_{\tilde{Q}_{L2}} = m_{\tilde{Q}_{L3}}$

Conclusion

- large parameter space leading to a light up-type squark (\approx light stop)
- $\tilde{u}_1 \rightarrow (c, u)\tilde{\chi}_1^0$ and $\tilde{u}_1 \rightarrow \tilde{\chi}_1^0 d_i \bar{f} f'$ for $\Delta m < m_W$
- two-body decay strongly dependent on the size of flavor mixing
- four-body decay nearly independent of flavor changing effects
- searches required in both the two-body and the four-body channel
- complete results and further details can be found in arXiv:1408.4662

Thanks for listening!

Conclusion

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Thanks for listening!

$$A_t \in [1000, 2000] \text{ GeV}$$

$$m_{\tilde{U}_3} \in [300, 600] \text{ GeV}$$

$$m_{\tilde{Q}_{L3}} \in [1000, 1500] \text{ GeV}$$

$$\tan \beta \in [1, 15]$$

$$M_1 \in [75, 500] \text{ GeV}$$

$$m_A \in [150, 1000] \text{ GeV}$$

All other parameters fixed:

$$M_2 = 650 \text{ GeV}$$

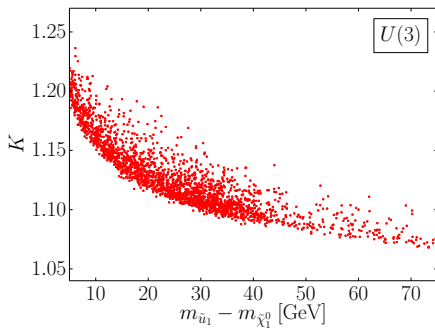
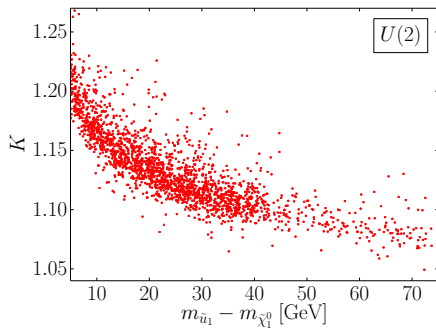
$$M_3 = 1530 \text{ GeV}$$

$$\mu = 900 \text{ GeV}$$

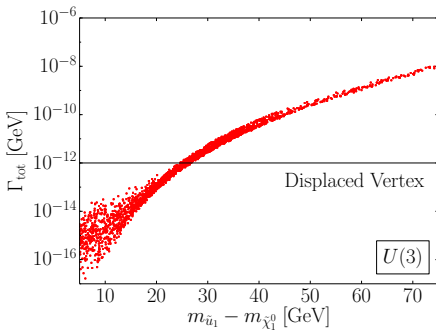
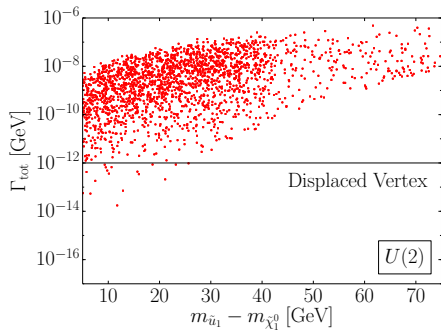
$$m_{\text{Sleptons}} = 1000 \text{ GeV}$$

$$m_{\text{Squarks}} = 1500 \text{ GeV}$$

$$A_d = A_l = 0$$

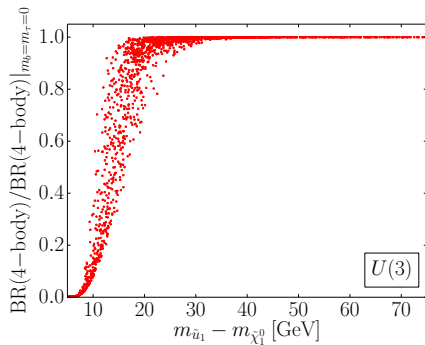
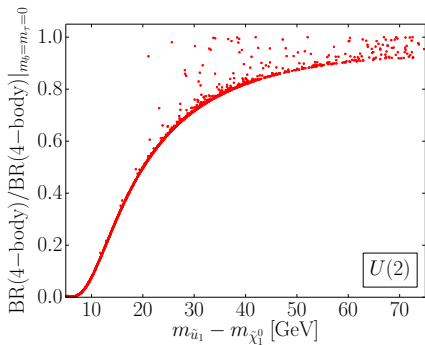


$$K = \frac{\Gamma_{\text{NLO}}}{\Gamma_{\text{LO}}}$$



$$\Gamma_{\text{tot}} = \Gamma_{\text{2-body}} + \Gamma_{\text{4-body}}$$

Backup: Effect of Massive Final State Particles



Definition:

[Djouadi *et al* '99]

- General MSSM with R -parity conservation and real parameters
- Minimal Flavor Violation, soft SUSY breaking masses and trilinear couplings are diagonal in flavor space
- trilinear couplings for the first two generations of sfermions can be neglected
- soft SUSY breaking masses for the first two generations of sfermions coincide

$$\begin{aligned} \begin{pmatrix} \tilde{f}_1 \\ \tilde{f}_2 \end{pmatrix} &= \begin{pmatrix} F_{11} & F_{12} \\ F_{21} & F_{22} \end{pmatrix} \begin{pmatrix} \tilde{f}_L \\ \tilde{f}_R \end{pmatrix} \\ \rightarrow \begin{pmatrix} \tilde{f}_1 \\ \tilde{f}_2 \\ \tilde{f}_3 \\ \tilde{f}_4 \\ \tilde{f}_5 \\ \tilde{f}_6 \end{pmatrix} &= \begin{pmatrix} F_{11} & \dots & \dots & \dots & F_{16} \\ \vdots & \ddots & & & \vdots \\ \vdots & & \ddots & & \vdots \\ \vdots & & & \ddots & \vdots \\ \vdots & & & & \ddots \\ F_{61} & \dots & \dots & \dots & F_{66} \end{pmatrix} \begin{pmatrix} \tilde{f}_{1L} \\ \tilde{f}_{2L} \\ \tilde{f}_{3L} \\ \tilde{f}_{1R} \\ \tilde{f}_{2R} \\ \tilde{f}_{3R} \end{pmatrix} \end{aligned}$$

Backup: Lagrangian for the Squark-Quark-Neutralino Interaction

$$\mathcal{L}_{\bar{u}\bar{u}\tilde{\chi}^0,L} = \underbrace{Q_{1il}^L \bar{u}_i^{(0)} \tilde{u}_{iR}^{(0)} \mathcal{P}_L \tilde{\chi}_l^0}_{\mathcal{L}_1} + \underbrace{Q_{2l}^L \bar{u}_i^{(0)} m_{ij}^{\dagger(0)} \tilde{u}_{jL}^{(0)} \mathcal{P}_L \tilde{\chi}_l^0}_{\mathcal{L}_2}$$

$$Q_{1il}^L := -g e_{Rl}^{u_i} = g\sqrt{2} Q_{u_i} t_W Z_{1l} \quad \text{and} \quad Q_{2l}^L := -\frac{g Z_{1l4}}{\sqrt{2} m_W s_\beta}$$

$$Q_{1il}^R := -g e_{Ll}^{u_i} = -g\sqrt{2} [Z_{1l} t_W (Q_{u_i} - I_{u_i}^3) + Z_{1l2} I_{u_i}^3] \quad \text{and} \quad Q_{2l}^R := -\frac{g Z_{1l4}}{\sqrt{2} m_W s_\beta}$$