Karlsruhe Institute of Technology

Institute for Theoretical Particle Physics

Top Flavoured Dark Matter in Dark Minimal Flavour Violation

GK Plenary Workshop, Freudenstadt

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KARLSRUHE INSTITUTE OF TECHNOLOGY

Outline

Simplified Models

- Presence of dark matter ($\Omega_{DM} \approx 27\%$) demands extension of Standard Model (SM).
  $\rightarrow$ What can we do to find the right extension?
- One extreme: full theory extension of SM (e.g. SUSY).
- Other extreme: effective field theory (EFT) approach.
- The middle way: simplified models.
- Advantage of simplified models: study specific interactions with limited number of parameters.

The Flavour Gate to Dark Matter

Assume an analogy to the SM fermions $\rightarrow$ dark flavour triplet χ_i.

The Flavour Gate to Dark Matter

Assume an analogy to the SM fermions $\to$ dark flavour triplet χ_i.

Flavoured dark matter coupling to SM right-handed up quark triplet:

$$\mathcal{L}_{\text{NP,int}} = -\lambda_{ij}\bar{u}_{Ri}\chi_j\phi + h.c.$$

- DM flavour triplet χ_j, Dirac fermion, SM gauge singlet.
- Heavy scalar mediator ϕ, carrying colour and hypercharge.
- Lagrangian has unbroken $\mathbb{Z}_3$ symmetry and hence yields stability of DM χ (for $m_\phi > m_\chi$).

Dark Minimal Flavour Violation

[Agrawal, Blanke, Gemmler '14]

Flavour symmetry

$$U(3)_u \times U(3)_d \times U(3)_q \times U(3)_\chi$$

is only broken by SM Yukawa couplings and the DM-quark coupling λ_{ij} (Dark Minimal Flavour Violation).

$\Rightarrow$ only DM mass splitting originates from RG running:

$$m_{ij} = m_\chi(\mathbb{1} + \eta\lambda^\dagger\lambda + ...)_{ij}.$$

- η depends on the full theory $\rightarrow$ has to be a parameter of the simplified model.
- Flavour with lowest mass is our DM candidate.
  $\rightarrow$ we choose the "top-flavour". [Kilic, Klimek, Yu '15]

Parametrization of DM-Quark Coupling Matrix

After using all the symmetries at our disposal, λ has 9 parameters left and can be parametrized as:

$$\boxed{\lambda = U_{23}^{\lambda} U_{13}^{\lambda} U_{12}^{\lambda} D_{\lambda}}$$

- D_λ is a real diagonal matrix $D_\lambda = \mathrm{diag}(D_{\lambda,11}, D_{\lambda,22}, D_{\lambda,33})$.
- U_{ij}^{λ} are unitary matrizes with mixing angles Θ_{ij} and phases δ_{ij}.

$\Rightarrow$ new source of flavour and CP violation

Constraints from SUSY Searches at LHC

Constraints from SUSY searches ($t\bar{t}$ or dijet final states) [ATLAS collaboration '14]

Study $pp \to \phi\bar{\phi} \to q\bar{q}\chi\bar{\chi}$

- Production either through $g\phi\bar{\phi}$ or NP interaction (coupling-dependent).
- Decay either to top or jet ($+\not{E}_T$).

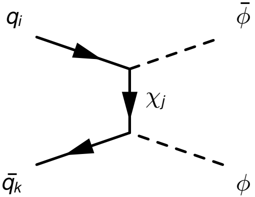

Figure : NP interaction production channel.

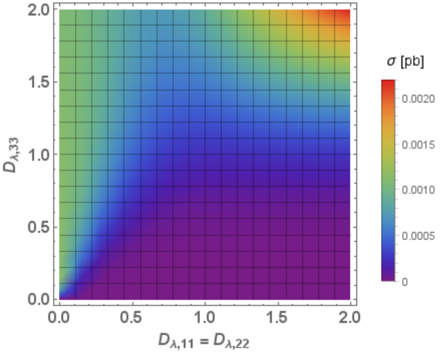

Figure : Cross section for $t\bar{t}$ final state, mixing angles set to zero.

Constraints from SUSY Searches at LHC

[ATLAS collaboration '14]

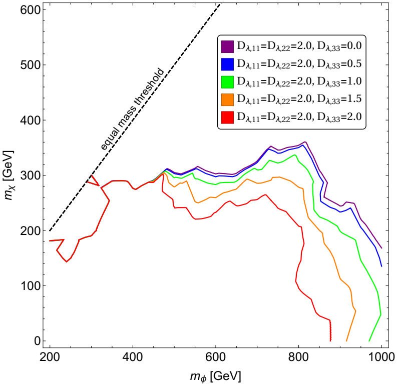

Figure : Exclusion plot for dijet final state, mixing angles set to zero.

- The phenomenologically interesting region is $m_\chi \leq 1$ TeV.
- Too large couplings $D_{\lambda,ii}$ would exclude nearly all of parameter space.
- Most serious constraints are given by searches for dijet final state.

$\Rightarrow$ Safe parameter space:

$$m_\phi \geq 850 \text{ GeV}$$

$$2.0 \geq D_{\lambda,33} > D_{\lambda,22}, D_{\lambda,11}$$

$\Rightarrow$ Also save with mixings allowed.

Flavour Constraints from Neutral Meson Mixing

[UTfit collaboration '14]

- No mesons with top quark are possible, the only constraints come from D mesons.
  $\Rightarrow$ not too strong

- The NP contribution has to be smaller than experimental bounds.
  $\Rightarrow$ constraints on mixing angles, mostly Θ_{12}

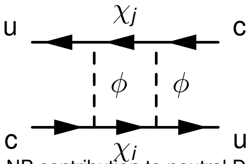

Figure : NP contribution to neutral D meson mixing.

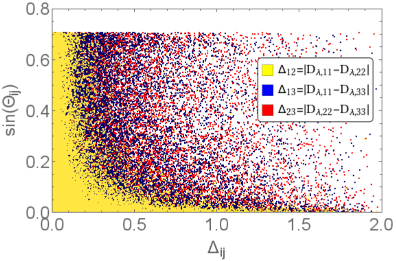

Figure : Valid mixing angles for different coupling splittings. $m_\phi = 850$ GeV, $m_\chi = 250$ GeV.

DM Constraints from Observed Relic Abundance

[Steigman, Dasgupta, Beacom ’12]

- Assume DM abundance as a thermal relic.
- Depending on mass splitting several freeze out scenarios are possible.
- If DM mass is below top mass several channels drop out.
  $\Rightarrow$ different impact on parameters
- Co-annihilation has to be just as large as to produce the correct relic density.
  $\Rightarrow$ cuts out valid area for $D_{\lambda,ii}$ depending on m_ϕ and m_χ
- Lower bounds on DM mass depending on mediator mass.
- Depending on η an upper DM bound arises in single flavour freeze out scenarios.

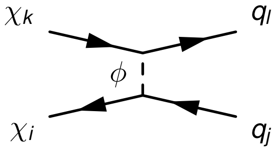

Figure : Coannihilation of DM flavours.

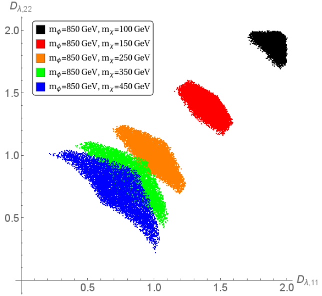

Figure : Valid points in quasi degenerate freeze out scenario.

DM Bounds from Direct Detection Experiments

[LUX collaboration '16]

- Many contributions to total WIMP-nucleon cross section, only Z-penguin with neutron is negative.
  $\Rightarrow$ saves the day
- Tree level and neutron Z-penguin have to nearly cancel each other.
  $\Rightarrow$ serious constraints on Θ_{13}
- For too large couplings the cancellation is no longer possible $\rightarrow$ excluded.
- Top flavoured DM is the natural choice.

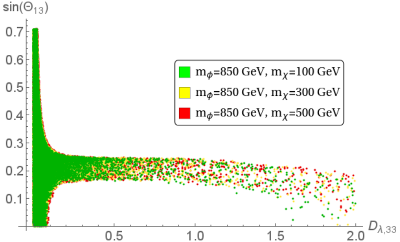

Figure : Valid mixing angle Θ_{13} vs $D_{\lambda,33}$.

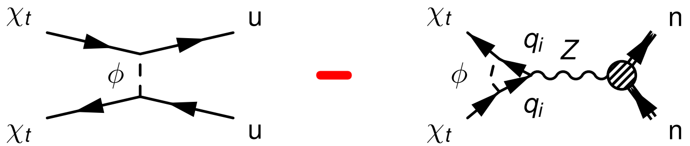

Figure : Cancellation of tree level and neutron Z-penguin contributions (symbolic).

Combined Analysis of Constraints

- A combination of relic abundance and direct detection constraints confine Θ_{13} to a narrow interval around the "perfect" cancellation point.
- The lower and upper bounds on the DM mass become more serious, since the parameters do not only have to fulfill relic abundance constraints.
- The combined analysis clearly prefers top flavoured DM.

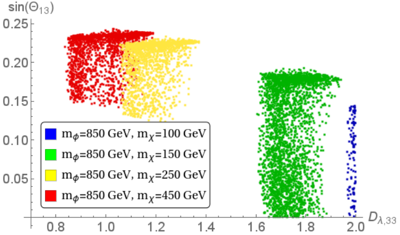

Figure : Valid points in Θ_{13}-$D_{\lambda,33}$-plane (QDF).

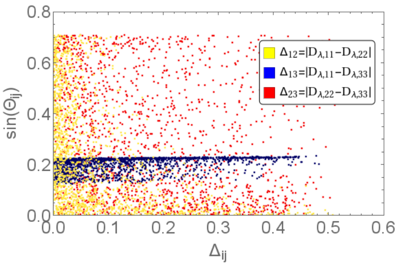

Figure : Valid points for $m_\phi = 850$ GeV and $m_\chi = 250$ GeV (QDF).

Recap

- A simplified model of flavoured DM coupled to SM right-handed up quark triplet. Coupling is general following the concept of DMFV.
- Assuming $m_\chi < 1$ TeV (phenomenologically interesting area).
- With this mass the RA constraints demand high $D_{\lambda,ii}$ for high mediator mass m_ϕ.
- High couplings prevent the necessary cancellation in WIMP-nucleon cross section. $\Rightarrow$ Mediator mass can not be too large if $m_\chi < 1$ TeV.
- Collider constraints limit couplings for a reasonable m_ϕ (NP production).
- Constraints from dijet searches prefer $D_{\lambda,33} \geq D_{\lambda,22}, D_{\lambda,11}$.
- Direct detection constraints prefer top flavoured DM.
- In combination with the limits on couplings, the RA constraints produce a lower bound for the DM mass (depending on m_ϕ).
- In SFF the splitting conditions in combination with RA constraints also establishes an upper bound on m_χ (depending on m_ϕ and η).

Conclusion and Outlook

- All kinds of different constraints $\rightarrow$ multitude of effects and interesting interplay.
- Especially interesting effect on mixing angle θ_{13} due to DD and RA constraints.
  $\Rightarrow$ Future measurements of direct detection experiments can potentially exclude a large class of models.
- Simplified models are powerful tool to study diversity of constraints.
- Going beyond Minimal Flavour Violation is worth the effort.
  $\rightarrow$ Dark Minimal Flavour Violation as guidance.

The End

Thank you!

The End

Thank you!

Questions?

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Constraints from SUSY Searches at LHC

[ATLAS collaboration '14]

- Study the process $pp \to \phi\bar{\phi} \to q\bar{q}\chi\bar{\chi}$.
- Depending on decay product of ϕ we detect either a top signature or a jet ($+\not{E}_T$).
- Inspiration from SUSY searches at LHC
  $\Rightarrow$ Upper bounds on CS of both $t\bar{t}$ and dijet signals.

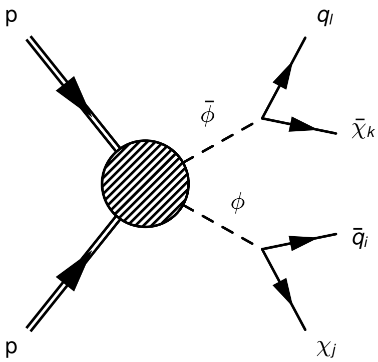
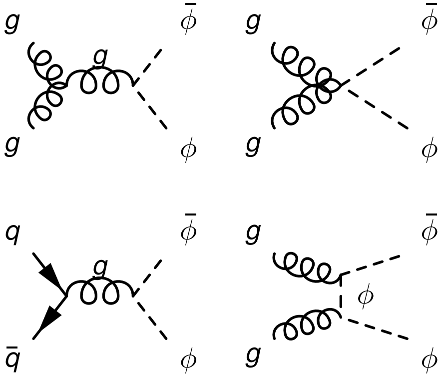


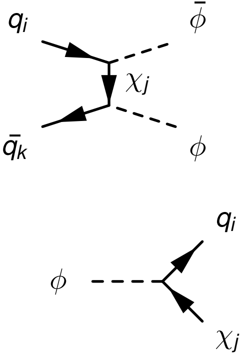
Figure : Studied LHC DM production processes.

Constraints from SUSY Searches at LHC

Involved QCD processes



Involved NP processes



Constraints from $t\bar{t}+\not{E}_T$ Searches at LHC

- $D_{\lambda,33}$ increased
  $\rightarrow$ BR of decay goes up.
- $D_{\lambda,11}$, $D_{\lambda,22}$ increased
  $\rightarrow$ BR of decay goes down.

- **BUT**: For high $D_{\lambda,11} = D_{\lambda,22}$ we observe increasing excluded areas.

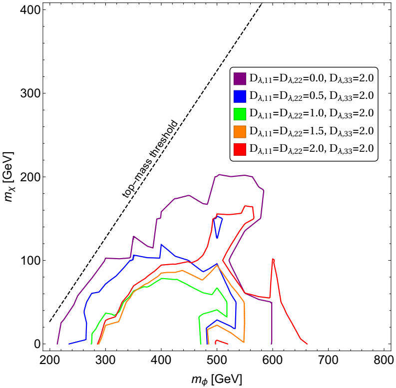


Figure : Exclusion plot for $t\bar{t}$ final state, mixing angles set to zero.

Constraints from SUSY Searches at LHC

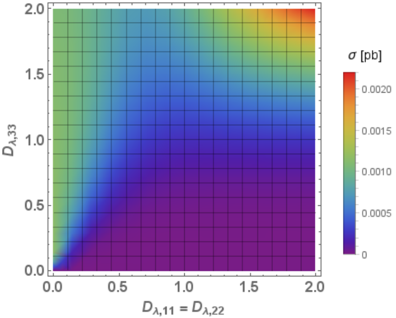


Figure : Cross section of $t\bar{t}$ final state for $m_\phi = 850$ GeV and $m_\chi = 50$ GeV, mixing angles set to zero.

Explanation: NP production

- Major contribution to total production (for high $D_{\lambda,11}$, $D_{\lambda,22}$)
- This effect can make up for drop in BR
- $D_{\lambda,33}$ not relevant, since the protons do not contain top
- Very high couplings can lead to serious exclusion areas.

Constraints from dijet + $\not{E}_T$ Searches at LHC

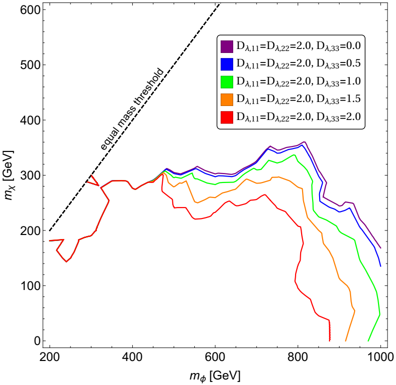


Figure : Exclusion plot for dijet final state, mixing angles set to zero.

- Stronger exclusion bounds on model.
- The phenomenologically interesting region is $m_\chi \leq 1$ TeV.
- Too large couplings $D_{\lambda,ii}$ would exclude nearly all of parameter space.
- Most serious constraints come from dijet final state.

$\Rightarrow$ Safe parameter space:

$$m_\phi \geq 850 \text{ GeV}$$

$$2.0 \geq D_{\lambda,33} \geq D_{\lambda,22}, D_{\lambda,11}$$

Influence of Mixing Angles on LHC production

- Mixing angles shift influences between couplings $D_{\lambda,ii}$.
  $\Rightarrow$ For big splitting in the couplings, mixing angles can cause big shifts in cross sections.
- For our choice of m_ϕ bounds from $t\bar{t}$ final state cause no constraints.
- Worst allowed case for dijet final state, in our safe parameter space, is $D_{\lambda,11} = D_{\lambda,22} = D_{\lambda,33} = 2.0$
  $\Rightarrow$ Unchanged by mixing angles.

$\Rightarrow$ Mixing angles can cause no problem with this choice of safe parameter space.

Flavour Constraints from Neutral Meson Mixing

[UTfit collaboration '14]

- No mesons with top quark are possible, the only constraints come from D mesons.
  $\Rightarrow$ not too strong
- The NP contribution has to be smaller than experimental bounds.

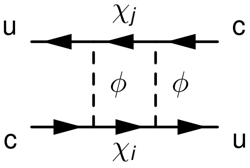


Figure : NP contr. to neutral D meson mixing.

$$
\begin{aligned}
M_{12}^{D,NP} &= \frac{1}{2m_D} \langle \bar{D}^0 | \mathcal{H}_{eff}^{\Delta C=2,new} | D^0 \rangle^* \\
&= \frac{1}{384\pi^2 m_\phi^2} \sum_{i,j} \lambda_{uj}^* \lambda_{cj} \lambda_{ui}^* \lambda_{ci} \cdot L(x_i, x_j) \cdot \eta_D \cdot m_D f_D^2 \hat{B}_D .
\end{aligned}
$$

Flavour Constraints from Neutral Meson Mixing

$$\left((\lambda\lambda^\dagger)_{cu}\right)^2 = \left((U_\lambda D_\lambda D_\lambda^\dagger U_\lambda^\dagger)_{cu}\right)^2$$

- For degeneracy $D_{\lambda,11} = D_{\lambda,22} = D_{\lambda,33}$ the mixing matrices U_{ij}^λ will drop out.
- The higher the splitting $\Delta_{ij} = D_{\lambda,ii} - D_{\lambda,jj}$, the more we will see the constraints on the mixing angle θ_{ij}.

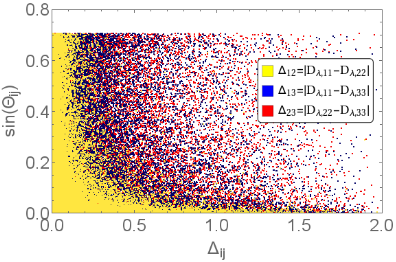


Figure : Valid mixing angles for different coupling splittings. $m_\phi = 850$ GeV and $m_\chi = 250$ GeV.

$\Rightarrow$ Most significant constraints on θ_{12}, other mixings nearly unconstrained.

DM Constraints from Observed Relic Abundance

[Steigman, Dasgupta, Beacom ’12]

- Assume DM abundance as a thermal relic, $T_f \propto \frac{m_\chi}{20}$
- Coannihilation CS has to be just large enough to produce the correct relic density (we allow for a 10% tolerance interval):

$$\langle \sigma v \rangle_{eff,exp} = 2.2 \times 10^{-26} \mathrm{cm}^3/\mathrm{s}.$$

$\Rightarrow$ cuts out valid area for $D_{\lambda,ii}$ depending on m_ϕ and m_χ

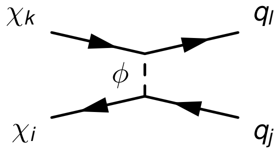


Figure : Coannihilation of DM flavours.

$$\langle \sigma v \rangle_{eff} = \frac{1}{9} \times \frac{3}{256\pi} \sum_{i,j=1,2,3} \sum_{k,l=u,c,t} \lambda_{ki}\lambda^*_{ki}\lambda_{lj}\lambda^*_{lj} \frac{\sqrt{\left(4m_\chi^2 - (m_k - m_l)^2\right)\left(4m_\chi^2 - (m_k + m_l)^2\right)}}{\left(m_\phi^2 + m_\chi^2 - \frac{m_k^2}{2} - \frac{m_l^2}{2}\right)^2}.$$

DM Constraints from Observed Relic Abundance

- Depending on the mass splitting of the different DM flavours several freeze out scenarios are possible.

$$m_{ij} = m_\chi(1 + \eta(D_{\lambda,ii})^2 + \ldots)\delta_{ij}.$$

- For a DM mass below the top quark mass this decay channel drops out.

  $\Rightarrow$ CS formula and hence impact on parameters can be quite different

- Extreme example: only χ_t present at freeze out with DM mass below top mass threshold:

$$\langle \sigma v \rangle_{eff} = \frac{3}{256\pi} \sum_{k,l=u,c} \lambda_{k3}\lambda_{k3}^*\lambda_{l3}\lambda_{l3}^* \frac{4m_\chi^2}{\left(m_\phi^2 + m_\chi^2\right)^2}.$$

Quasi Degenerate Freeze Out (QDF) Szenario

- All DM flavours are present at the freeze out.
- We require the mass splitting to be less than 1% (significantly smaller than T_f) for this to happen.
- η is free parameter $\rightarrow$ choose it favourable: -0.01.
- This guarantees top flavoured DM (see direct detection section for motivation).
- Constraint cuts out valid area for $D_{\lambda,ii}$ depending on m_ϕ and m_χ.
- Lower bound on m_χ due to upper limits for $D_{\lambda,ii}$, depending on m_ϕ.

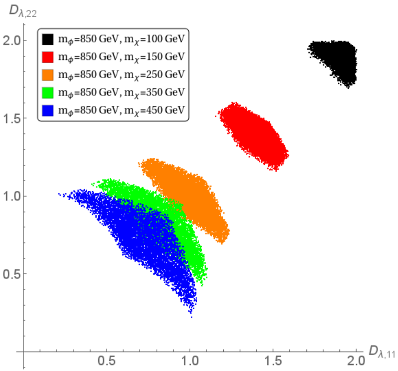


Figure : Valid points in quasi degenerate freeze out scenario.

Single Flavour Freeze Out (SFF) Szenario

- Only m_χ present at freeze out.
- We require the mass splitting to be more than 10% (significantly bigger than T_f) for this to happen.
- η is free parameter $\rightarrow$ choose it favourable: -0.075.
- This guarantees top flavoured DM (see direct detection section for motivation).
- Constraint cuts out valid area of parameters depending on m_ϕ and m_χ, with significant effect on mixing angles.
- In addition to lower bound, we also find an upper bound on m_χ due to upper and lower (from mass splitting condition) limits for $D_{\lambda,ii}$, depending on m_ϕ.

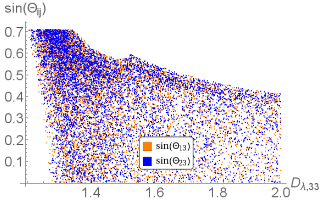


Figure : Valid points in single flavour freeze out scenario for $m_\phi = 850$ GeV and $m_\chi = 210$ GeV.

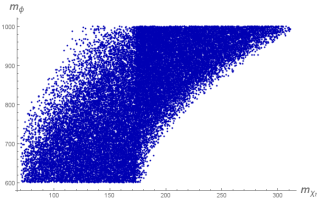
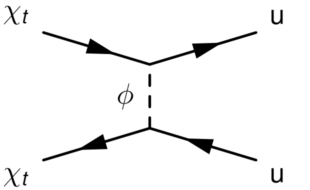


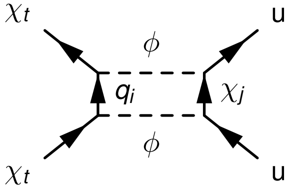
Figure : Mass bounds in single flavour freeze out scenario.

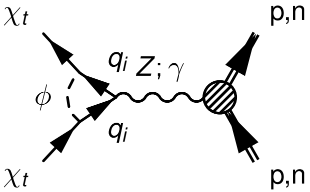
DM Bounds from Direct Detection Experiments

Many contributions to total WIMP-nucleon cross section:

$$\sigma_n^{SI} = \frac{\mu_n^2}{\pi A^2} |Zf_p + (A - Z) f_n|^2.$$







References

DM Bounds from Direct Detection Experiments

[LUX collaboration '15]

- All contributions have to combine to a WIMP-nucleon cross section below the LUX bounds.
- All contributions are positive, only the Z-penguin with the neutron is negative $\Rightarrow$ saves the day.
- Largest contribution comes from tree level process. Largest negative term is hence interference term of tree level and neutron Z-penguin.
- Most important terms, have to nearly cancel each other:

$$A_{\mathcal{I}} \cdot D^4_{\lambda,33} \cdot sin(\theta_{13})^4 - A_{\mathcal{I}\mathcal{I}} \cdot D^4_{\lambda,33} \cdot sin(\theta_{13})^2 \cdot cos(\theta_{13})^2 \cdot cos(\theta_{23})^2$$

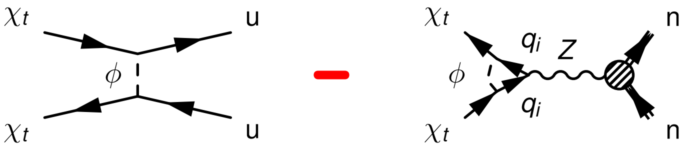


Figure : Cancellation of tree level and neutron Z-penguin contributions (symbolic).

DM Bounds from Direct Detection Experiments

- Tree level and neutron Z-penguin have to nearly cancel each other.
  $\Rightarrow$ serious constraints on θ_{13}
- For higher couplings the cancellation gets more complicated.
- For too large couplings the cancellation is no longer possible at all
  $\rightarrow$ excluded.
- Top-flavoured DM is the natural choice:
  $\Rightarrow$ Tree level contribution small
  $\Rightarrow$ Neutron Z-penguin contribution large.

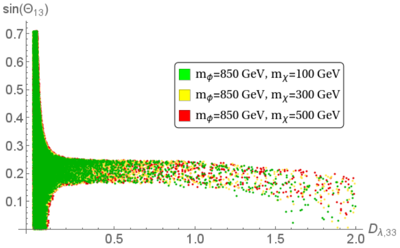


Figure : Valid mixing angle Θ_{13} vs $D_{\lambda,33}$.

Combined Analysis of Constraints (QDF)

Combined application of both flavour, relic abundance and direct detection constraint in quasi degenerate freeze out scenario.

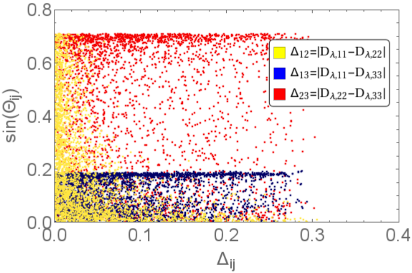


Figure : Valid points for $m_\phi = 850$ GeV and $m_\chi = 150$ GeV (QDF).

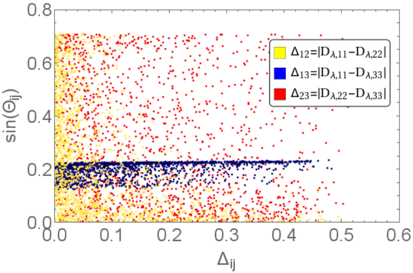


Figure : Valid points for $m_\phi = 850$ GeV and $m_\chi = 250$ GeV (QDF).

Combined Analysis of Constraints (QDF)

- A combination of relic abundance and direct detection constraints confine θ_{13} to a narrow interval.
- The bounds on the DM mass become more serious, since the parameters do not only have to fulfill relic abundance constraints.
- The combined analysis clearly prefers top flavoured DM.

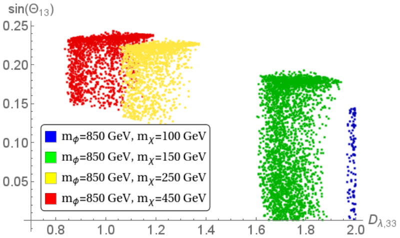


Figure : Valid points in θ_{13}-$D_{\lambda,33}$-plane (QDF).

Combined Analysis of Constraints (SFF)

Combined application of both flavour, relic abundance and direct detection constraint in single flavour freeze out scenario.

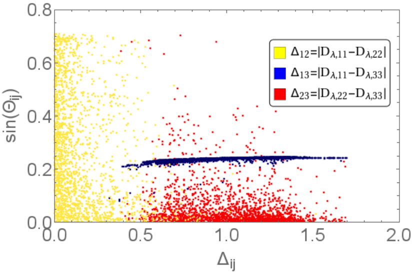


Figure : Valid points for $m_\phi = 850$ GeV and $m_\chi = 225$ GeV (SFF).

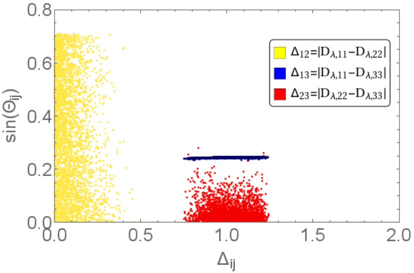


Figure : Valid points for $m_\phi = 850$ GeV and $m_\chi = 250$ GeV (SFF).

Combined Analysis of Constraints (SFF)

- A combination of relic abundance and direct detection constraints confine θ_{13} to a narrow interval (even more serious than in QDF).
- Especially in SFF the combination of all constraints extremely limits the chance of finding a valid configuration of all parameters for $m_{\chi_t} \leq m_{top}$.
- The combined analysis clearly prefers top flavoured DM.

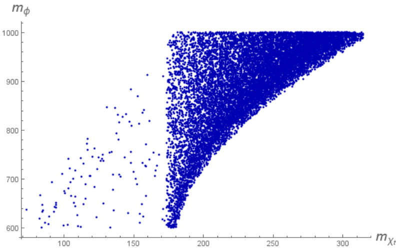


Figure : Valid points in mass plot for combined constraints (SFF).

References

Backup Material 1

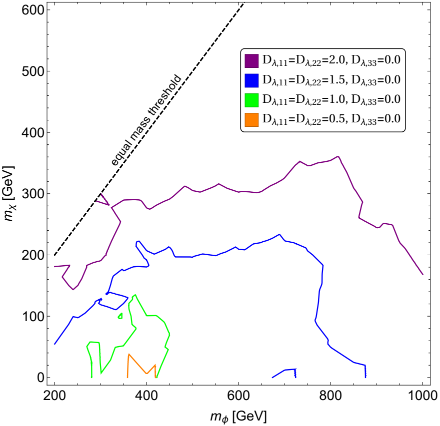


Figure : Exclusion plots for dijet final state for various couplings, mixing angles set to zero.

Backup Material 2

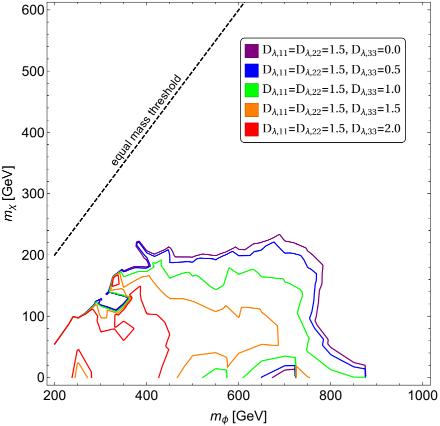


Figure : Exclusion plots for dijet final state for various couplings, mixing angles set to zero.

Backup Material 3

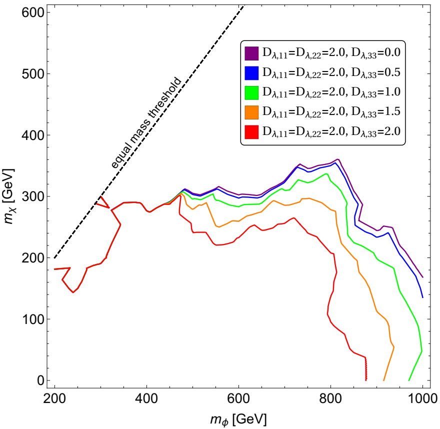


Figure : Exclusion plots for dijet final state for various couplings, mixing angles set to zero.

Backup Material 4

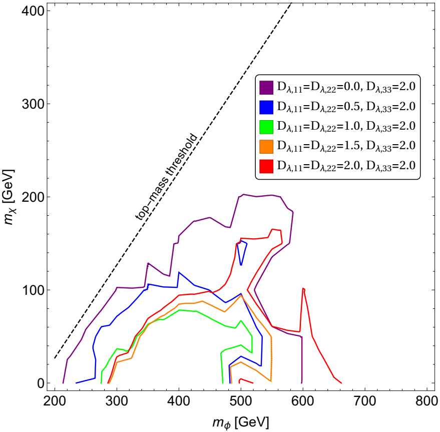


Figure : Exclusion plots for $t\bar{t}$ final state for various couplings, mixing angles set to zero.

Backup Material 5

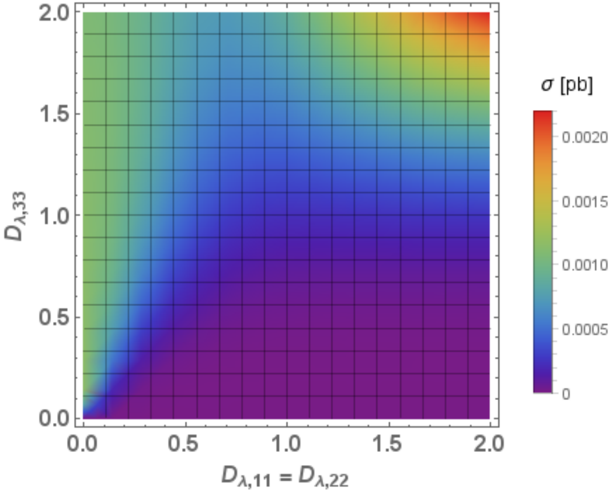


Figure : Cross section for $t\bar{t}$ final state, mixing angles set to zero.

Backup Material 6

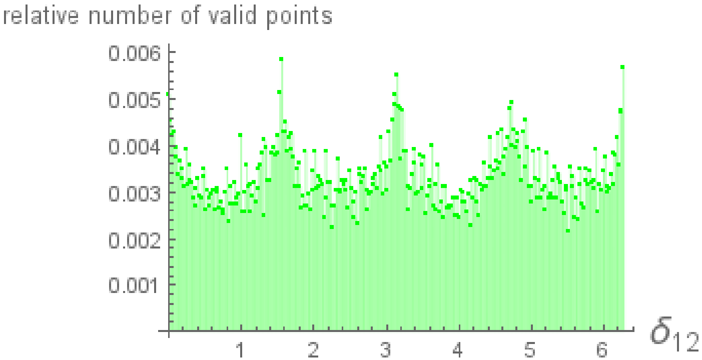


Figure : Impact of flavour constraints on Θ_{12}.

Backup Material 7

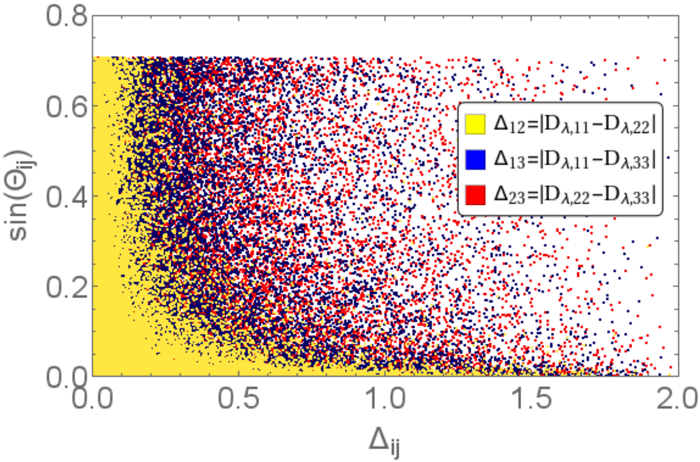


Figure : Valid mixing angles for different coupling splittings. $m_\phi = 850$ GeV and $m_\chi = 250$ GeV.

Backup Material 8

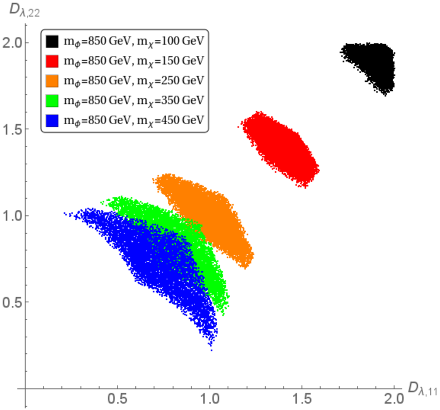


Figure : Valid points in quasi degenerate freeze out scenario in $D_{\lambda,11} - D_{\lambda,22}$-plane for various DM masses.

Backup Material 9

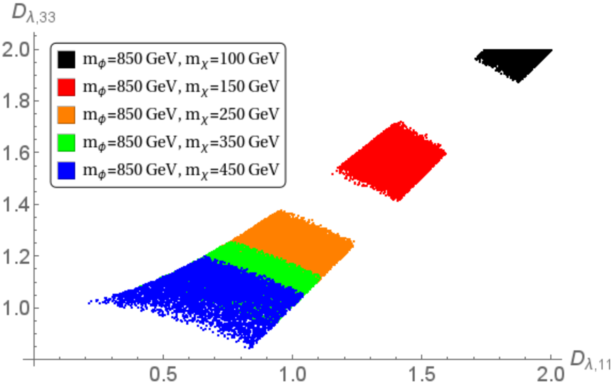


Figure : Valid points in quasi degenerate freeze out scenario in $D_{\lambda,11} - D_{\lambda,33}$-plane for various DM masses.

Backup Material 10

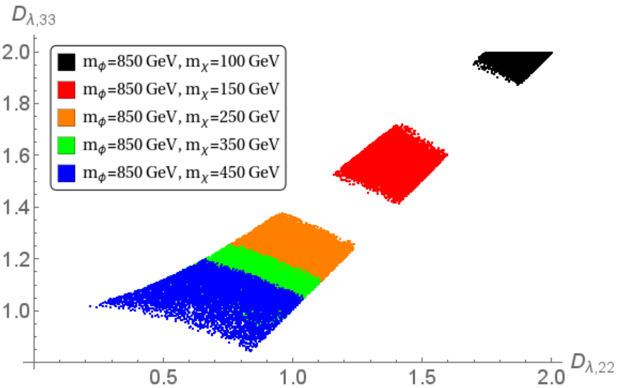


Figure : Valid points in quasi degenerate freeze out scenario in $D_{\lambda,22} - D_{\lambda,33}$-plane for various DM masses.

Backup Material 11

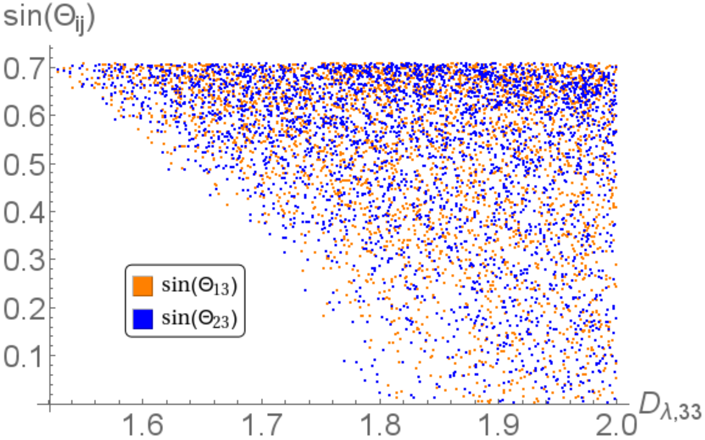


Figure : Valid points in single flavour freeze out scenario in $D_{\lambda,33} - \sin(\Theta_{ij})$-plane for $m_\phi = 850$ GeV and $m_\chi = 150$ GeV.

Backup Material 12

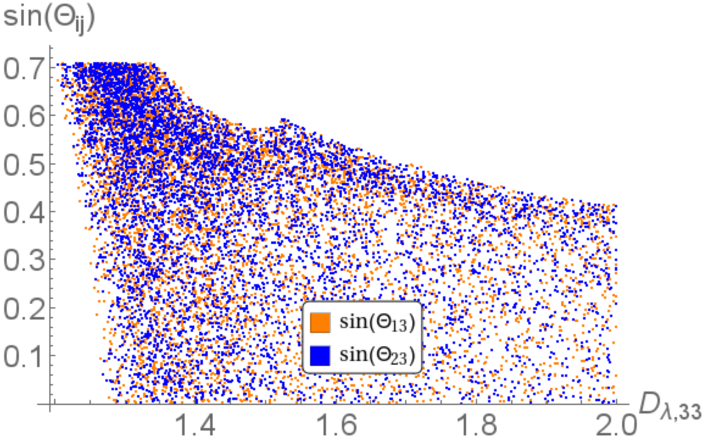


Figure : Valid points in single flavour freeze out scenario in $D_{\lambda,33} - \sin(\Theta_{ij})$-plane for $m_\phi = 850$ GeV and $m_\chi = 210$ GeV.

Backup Material 13

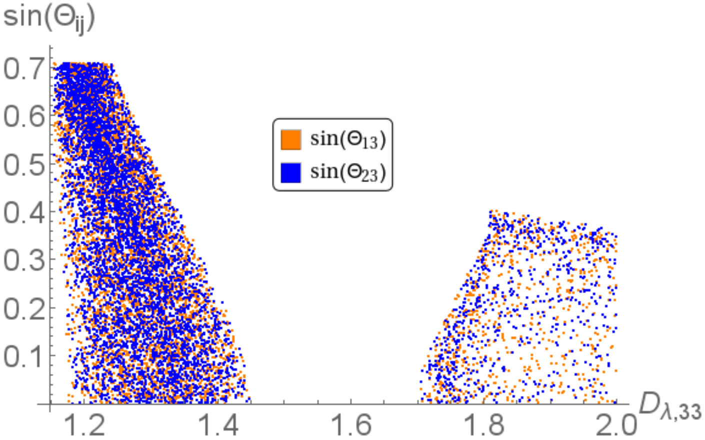


Figure : Valid points in single flavour freeze out scenario in $D_{\lambda,33} - \sin(\Theta_{ij})$-plane for $m_\phi = 850$ GeV and $m_\chi = 230$ GeV.

Backup Material 14

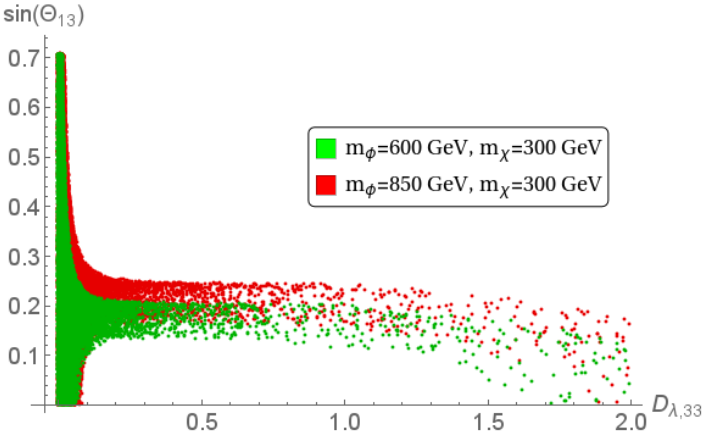


Figure : Valid points for LUX bounds in $D_{\lambda,33} - \sin(\Theta_{13})-$plane.

Backup Material 15

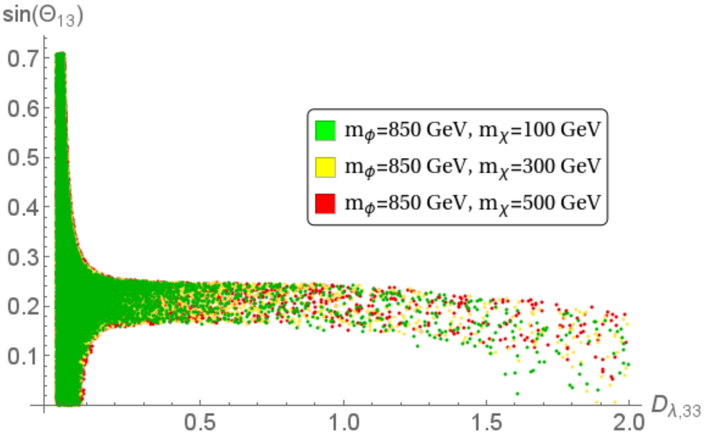


Figure : Valid points for LUX bounds in $D_{\lambda,33} - \sin(\Theta_{13})-$plane.

Backup Material 16

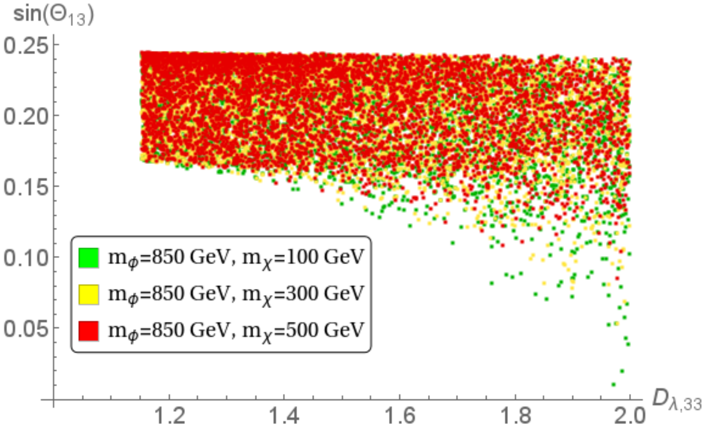


Figure : Valid points for LUX bounds in $D_{\lambda,33} - \sin(\Theta_{13})$−plane, with SFF splitting applied.

Backup Material 17

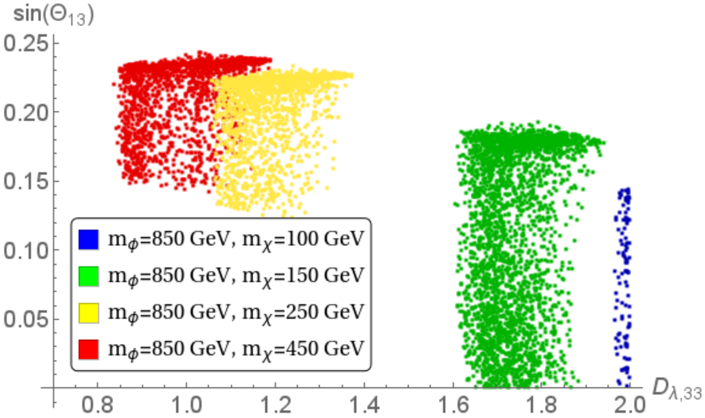


Figure : Valid points of combined analysis for quasi degenerate freeze out scenario in $D_{\lambda,33} - \sin(\Theta_{13})-$plane for different DM masses.

Backup Material 18

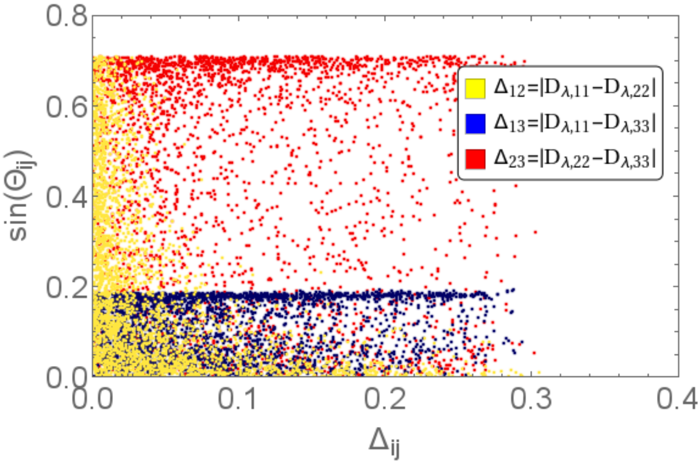


Figure : Valid mixing angles for different coupling splittings for quasi degenerate freeze out scenario. $m_\phi = 850$ GeV and $m_\chi = 150$ GeV.

Backup Material 19

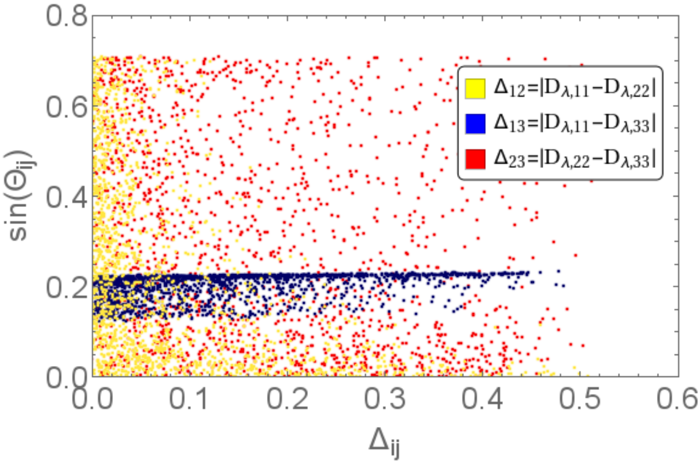


Figure : Valid mixing angles for different coupling splittings for quasi degenerate freeze out scenario. $m_\phi = 850$ GeV and $m_\chi = 250$ GeV.

Backup Material 20

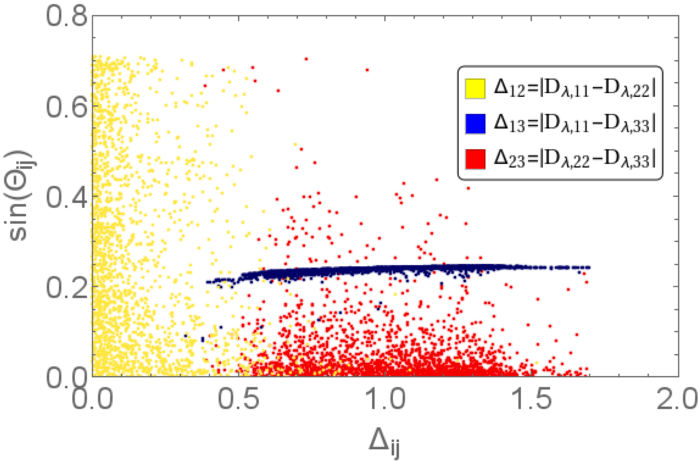


Figure : Valid mixing angles for different coupling splittings for single flavour freeze out scenario. $m_\phi = 850$ GeV and $m_\chi = 225$ GeV.

Backup Material 21

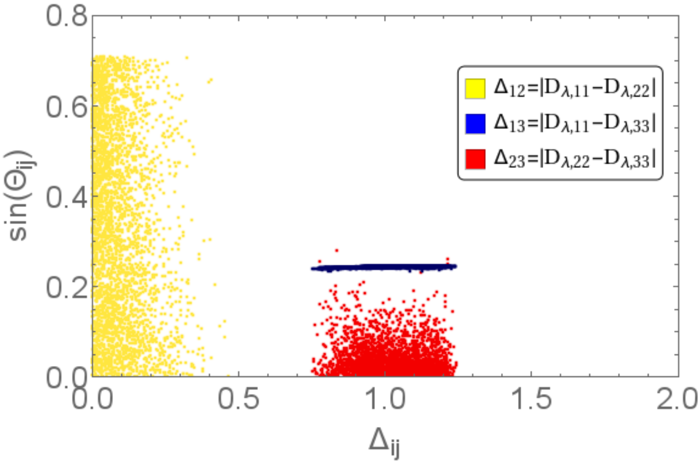


Figure : Valid mixing angles for different coupling splittings for single flavour freeze out scenario. $m_\phi = 850$ GeV and $m_\chi = 250$ GeV.